

We would now like to compare the moment of inertia for a rigid body.

Let's take an arbitrary rigid body about the center of mass.

So let's say the rigid body is rotating about this axis.

And what we'd like to compare that is to the moment of inertia, say, about a parallel axis that's also going through the rigid body.

Now, let's recall how we define moment of inertia.

We first choose a mass element  $dm$ .

What I'd like to show is if this object is rotating about this axis, then what is that object doing.

Well, that object is undergoing a little bit of circular motion.

And this distance here is what we call the perpendicular distance about that axis.

And let's indicate this is for our element  $dm$ .

So this perpendicular distance is what shows up in our definition for the center of mass moment of inertia about that axis-- it's the interval of  $dm r_{cm\ perp}$ .

Now again, quantity squared.

What is this distance?

This is the perpendicular distance from our  $dm$  and to the axis of rotation.

Imagine it's doing a circle and that's the radius of that circle.

So if we were to calculate the moment of inertia about another axis, then about this axis the perpendicular distance here that I'll write as  $r_{s\ perp}$ .

And you can see these perpendicular distances are not the same.

And the moment of inertia about that other axis is equal to the integral of  $dm r_{s\ perp}$  quantity squared.

Now, how do we relate those perpendicular distances?

Well, there's a couple of ways to do it.

And notice that the distance between these axes is given by  $d$ .

And I'm going to call this the distance  $r_{cm}$ .

Now, let's just call this-- the  $x$  direction-- I'll call that  $x$ .

So how do I relate these distances?

Well,  $d$  is a fixed distance.

And you can see from my diagram that  $r_{\text{perp}}$  is equal to  $d + r_{cm} \cos \theta$ .

And if I square this, I get  $d^2 + 2d r_{cm} \cos \theta + r_{cm}^2 \cos^2 \theta$ .

And that  $r_{cm}^2 \cos^2 \theta$  is precisely what we're calling perpendicular distance.

So when I put those into my moment of inertia  $I$ s, I get  $\int dm d^2 + 2d \int dm r_{cm} \cos \theta + \int dm r_{\text{perp}}^2$ .

Now, I'll separate this into three terms.

The first term is  $d^2 \int dm$ .

This is an integral over the body.

The second term is  $2d \int dm r_{cm} \cos \theta$  and I'm going to hold off on the interval, because the  $2d$  is the same for every piece--  $\int dm r_{cm} \cos \theta$ .

And the third piece is  $\int dm r_{\text{perp}}^2$  since  $r_{cm} \cos \theta$  is the  $r_{\text{perp}}$ , I'll write it as  $r_{\text{perp}}^2$ .

And you can see that this term is precisely the moment of inertia about the center of mass.

Now, what I'd like to focus on is this term,  $\int dm r_{cm} \cos \theta$  that appears in our integral expression.

Recall, that we define center of mass.

We had the condition that the sum of  $m_j r_{cmj}$  was 0.

Now for an integral relationship, this is  $\int dm r_{cm}$  equal to 0.

So when you sum up the position of every object with respect to the vector from the center of mass to your  $dm$  element, 0.

What does this say in terms of components?

In terms of components, each component separately vanishes so we have the condition that  $\sum m_i x_i = 0$ .

So that term is 0, which is precisely this term-- that's 0.

And so we can conclude that  $I = \sum m_i d_i^2$  now in this term, where  $d$  is the same piece for every object-- so we're just pulling out the total mass.

So it's  $m$  total  $d$ -squared.

And let's remind ourselves that  $d$  is the distance between the two parallel axes plus the moment of inertia about the center of mass.

And this result is called the parallel axis theorem.