# 18.303: Self-adjointness (reciprocity) and definiteness (positivity) in Green's functions 

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## 0 Review

Suppose we have some vector space $V$ of functions $u(\mathbf{x})$ on a domain $\Omega$, an inner product $\langle u, v\rangle$, and a linear operator $\hat{A}$. [More specifically, $V$ forms a Sobolev space, in that we require $\langle u, \hat{A} u\rangle$ to be finite.] $\hat{A}$ is self-adjoint if $\langle u, \hat{A} v\rangle=\langle\hat{A} u, v\rangle$ for all $u, v \in V$, in which case its eigenvalues $\lambda_{n}$ are real and its eigenfunctions $u_{n}(\mathbf{x})$ can be chosen orthonormal. $\hat{A}$ is positive definite (or semidefinite) if $\langle u, \hat{A} u\rangle>0($ or $\geq 0)$ for all $u \neq 0$, in which case its eigenvalues are $>0$ (or $\geq 0$ ); suppose that we order them as $0<\lambda_{1} \leq \lambda_{2} \leq \cdots$.

Suppose that $\hat{A}$ is positive definite, so that $N(\hat{A})=\{0\}$ and $\hat{A} u=f$ has a unique solution for all $f$ in some suitable space of functions $C(\hat{A})$. Then, for scalar-valued functions $u$ and $f$, we can typically write

$$
\begin{equation*}
u(\mathbf{x})=\hat{A}^{-1} f=\int_{\mathbf{x}^{\prime} \in \Omega} G\left(\mathbf{x}, \mathbf{x}^{\prime}\right) f\left(\mathbf{x}^{\prime}\right) \tag{1}
\end{equation*}
$$

in terms of a Green's function $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)$, where $\int_{\mathbf{x}^{\prime} \in \Omega}$ denotes integration over $\mathbf{x}^{\prime}$. In this note, we don't address how to find $G$, but instead ask what properties it must have from the self-adjointness and definiteness of $\hat{A}$. [This generalizes in a straightforward way to vector-valued $\mathbf{u}(\mathbf{x})$ and $\mathbf{f}(\mathbf{x})$, in which case $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)$ is matrix-valued.]

## 1 Self-adjointness of $\hat{A}^{-1}$ and reciprocity of $G$

We can show that $\left(\hat{A}^{-1}\right)^{*}=\left(\hat{A}^{*}\right)^{-1}$, from which it follows that if $\hat{A}=\hat{A}^{*}(\hat{A}$ is self-adjoint) then $\hat{A}^{-1}$ is also self-adjoint. In particular, consider $\hat{A}^{-1} \hat{A}=1:\langle u, v\rangle=\left\langle u, \hat{A}^{-1} \hat{A} v\right\rangle=\left\langle\left(\hat{A}^{-1}\right)^{*} u, \hat{A} v\right\rangle=$ $\left\langle\hat{A}^{*}\left(\hat{A}^{-1}\right)^{*} u, v\right\rangle$, hence $\hat{A}^{*}\left(\hat{A}^{-1}\right)^{*}=1$ and $\left(\hat{A}^{-1}\right)^{*}=\left(\hat{A}^{*}\right)^{-1}$. And of course, we already knew that the eigenvalues of $\hat{A}^{-1}$ are $\lambda_{n}^{-1}$ and the eigenfunctions are $u_{n}(\mathbf{x})$.

What are the consequences of self-adjointness for $G$ ? Suppose the $u$ are scalar functions, and that the inner product is of the form $\langle u, v\rangle=\int_{\Omega} w \bar{u} v$ for some weight $w(\mathbf{x})>0$. From the fact that $\left\langle u, \hat{A}^{-1} v\right\rangle=\left\langle\hat{A}^{-1} u, v\right\rangle$, substituting equation (1), we must therefore have:

$$
\begin{aligned}
\left\langle u, \hat{A}^{-1} v\right\rangle & =\iint_{\mathbf{x}, \mathbf{x}^{\prime} \in \Omega} w(\mathbf{x}) \overline{u(\mathbf{x})} G\left(\mathbf{x}, \mathbf{x}^{\prime}\right) v\left(\mathbf{x}^{\prime}\right) \\
& =\left\langle\hat{A}^{-1} u, v\right\rangle \\
& =\iint_{\mathbf{x}, \mathbf{x}^{\prime} \in \Omega} w(\mathbf{x}) \overline{G\left(\mathbf{x}, \mathbf{x}^{\prime}\right) u\left(\mathbf{x}^{\prime}\right)} v(\mathbf{x})=\iint_{\mathbf{x}, \mathbf{x}^{\prime} \in \Omega} w\left(\mathbf{x}^{\prime}\right) \overline{u(\mathbf{x}) G\left(\mathbf{x}^{\prime}, \mathbf{x}\right)} v\left(\mathbf{x}^{\prime}\right),
\end{aligned}
$$

where in the last step we have interchanged/relabeled $\mathbf{x} \leftrightarrow \mathbf{x}^{\prime}$. Since this must be true for all $u$ and $v$, it follows that

$$
w(\mathbf{x}) G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)=w\left(\mathbf{x}^{\prime}\right) \overline{G\left(\mathbf{x}^{\prime}, \mathbf{x}\right)}
$$

for all $\mathbf{x}, \mathbf{x}^{\prime}$. This property of $G$ (or its analogues in other systems) is sometimes called reciprocity. In the common case where $w=1$ and $\hat{A}$ and $G$ are real (so that the complex conjugation can be omitted), it says that the effect at $\mathbf{x}$ from a source at $\mathbf{x}^{\prime}$ is the same as the effect at $\mathbf{x}^{\prime}$ from a source at $\mathbf{x}$.

There are many interesting consequences of reciprocity. For example, its analogue in linear electrical circuits says that the current at one place created by a voltage at another is the same as if the locations of the current and voltage are swapped. Or, for antennas, the analogous theorem says that a given antenna works equally well as a transmitter or a receiver.

### 1.1 Example: $\hat{A}=-\frac{d^{2}}{d x^{2}}$ on $\Omega=[0, L]$

For this simple example (where $\hat{A}$ is self-adjoint under $\langle u, v\rangle=\int \bar{u} v$ ), with Dirichlet boundaries, we previously obtained a Green's function,

$$
G\left(x, x^{\prime}\right)= \begin{cases}\left(1-\frac{x^{\prime}}{L}\right) x & x<x^{\prime} \\ \left(1-\frac{x}{L}\right) x^{\prime} & x \geq x^{\prime}\end{cases}
$$

which obviously obeys the $G\left(x, x^{\prime}\right)=G\left(x^{\prime}, x\right)$ reciprocity relation.

## 2 Positive-definiteness of $\hat{A}^{-1}$ and positivity of $G$

Not only is $\hat{A}^{-1}$ self-adjoint, but since its eigenvalues are the inverses $\lambda_{n}^{-1}$ of the eigenvalues of $\hat{A}$, then if $\hat{A}$ is positive-definite $\left(\lambda_{n}>0\right)$ then $\hat{A}^{-1}$ is also positive-definite $\left(\lambda_{n}^{-1}>0\right)$. From another perspective, if $\hat{A} u=f$, then positive-definiteness of $\hat{A}$ means that $0<\langle u, \hat{A} u\rangle=\langle u, f\rangle=\left\langle\hat{A}^{-1} f, f\right\rangle=$ $\left\langle f, \hat{A}^{-1} f\right\rangle$ for $u \neq 0 \Leftrightarrow f \neq 0$, hence $\hat{A}^{-1}$ is positive-definite. (And if $\hat{A}$ is a PDE operator with an ascending sequence of unbounded eigenvalues, then the eigenvalues of $\hat{A}^{-1}$ are a descending sequence $\lambda_{1}^{-1}>\lambda_{2}^{-1}>\cdots>0$ that approaches 0 asymptotically from above. $\frac{1}{\text { ) }}$ )

If $\hat{A}$ is a real operator (real $u$ give real $\hat{A} u$ ), then $\hat{A}^{-1}$ should also be a real operator (real $f$ give real $u=\hat{A}^{-1} f$ ). Furthermore, under fairly general conditions for real positive-definite (elliptic) PDE operators $\hat{A}$, especially for second-derivative ("order 2") operators, then one can often show $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)>0$ (except of course for $\mathbf{x}$ or $\mathbf{x}^{\prime}$ at the boundaries, where $G$ vanishes for Dirichlet conditions). $\underset{\sim}{2}$ The analogous fact for matrices $A$ is that if $A$ is real-symmetric positive-definite and it has off-diagonal entries $\leq 0-$ like our $-\nabla^{2}$ second-derivative matrices (recall the $-1,2,-1$ sequences in the rows) and related finite-difference matrices - it is called a Stieltjes matrix, and such matrices can be shown to have inverses with nonnegative entries. $-\frac{3}{-}$

### 2.1 Example: $\hat{A}=-\nabla^{2}$ with $\left.u\right|_{\partial \Omega}=0$

Physically, the positive-definite problem $-\nabla^{2} u=f$ can be thought of as the displacement $u$ in response to an applied pressure $f$, where the Dirichlet boundary conditions correspond to a material pinned at the edges. The Green's function $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)$ is the limit of the displacement $u$ in response to a force concentrated at a single point $\mathbf{x}^{\prime}$. The Green's function $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)$ for some example points $\mathbf{x}^{\prime}$ is shown for a 1 d domain $\Omega=[0,1]$ in figure 1 (left) (a "stretched string"), and for a 2 d domain $\Omega=[-1,1] \times[-1,1]$ in figure 1(right) (a "square drum"). As expected, $G>0$ everywhere except at the edges where it is zero: the whole string/membrane moves in the positive/upwards direction in response to a positive/upwards force.

[^0]

Figure 1: Examples illustrating the positivity of the Green's function $G\left(\mathbf{x}, \mathbf{x}^{\prime}\right)$ for a positive-definite operator ( $-\nabla^{2}$ with Dirichlet boundaries). Left: a "stretched string" 1d domain $[0,1]$. Right: a "stretched square drum" 2 d domain $[-1,1] \times[-1,1]$.

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[^0]:    ${ }^{1}$ Such $\hat{A}^{-1}$ integral operators are typically what are called "compact" operators. Functional analysis books often prove diagonalizability (a "spectral theorem") for compact operators first and only later consider diagonalizability of PDE-like operators by viewing them as the inverses of compact operators.
    ${ }^{2}$ See, for example, "Characterization of positive reproducing kernels. Application to Green's functions," by N. Aronszajn and K. T. Smith [Am. J. Mathematics, vol. 79, pp. 611-622 (1957), http://www.jstor.org/stable/2372564]. However, as usual there are pathological counter-examples.
    ${ }^{3}$ There are many books with "nonnegative matrices" in their titles that cover this fact, usually as a special case of a more general class of something called "M matrices," but I haven't yet found an elementary presentation at an 18.06 level. Note that the diagonal entries of a positive-definite matrix $P$ are always positive, thanks to the fact that $P_{i i}=\mathbf{e}_{i}^{T} P \mathbf{e}_{i}>0$ where $\mathbf{e}_{i}$ is the unit vector in the $i$-th coordinate.

