## The decay of solutions of the heat equation, Campanato's lemma, and Morrey's Lemma

## 1 The decay of solutions of the heat equation

A few lectures ago we introduced the heat equation

$$\Delta u = u_t \tag{1}$$

for functions of both space and time. we will now bound the decay of the Dirichlet energy and the  $L_2$  norm (ie  $\int u^2$ ) of solutions. If u solves the heat equation on  $\Omega \subset \mathbb{R}^n$  with u = 0on the boundary then

$$\frac{d}{dt} \int_{\Omega} u^2 = 2 \int_{\Omega} u u_t = 2 \int_{\Omega} u \triangle u = -2 \int_{\Omega} |\nabla u|^2, \tag{2}$$

and

$$\frac{d}{dt} \int_{\Omega} |\nabla u|^2 = 2 \int_{\Omega} \nabla u \cdot \nabla u_t = -2 \int_{\Omega} u_t \Delta u = -2 \int_{\Omega} |\Delta u|^2.$$
(3)

Also note that the dirichlet energy is decreasing, so

$$-2\int_{\Omega} |\nabla u(\cdot, 0)|^2 \le \frac{d}{dt} \int_{\Omega} u^2 \le -2\int_{\Omega} |\nabla u(\cdot, t)|^2, \tag{4}$$

and

$$-2T \int_{\Omega} |\nabla u(\cdot, 0)|^2 \le \int_{\Omega} u^2(\cdot, T) - \int_{\Omega} u^2(\cdot, 0) \le -2T \int_{\Omega} |\nabla u(\cdot, T)|^2.$$
(5)

## 2 Campanato's Lemma

Before discussing Campanato's lemma we need the concept of Holder continuity.

**Definition** Let  $f : \mathbb{R}^n \to \mathbb{R}$  be a continuous function. f is  $\alpha$  continuous at  $x_0$  if there are constants k and  $0 < \alpha \leq 1$  such that

$$f(x_0 + x) - f(x_0) \le k|x|^{\alpha}$$
 (6)

for small  $x_0$ . If f is everywhere Holder continuous for some constant  $\alpha$  then we say that  $f \in C^{\alpha}$ .

Now take f continuous on  $\mathbb{R}^n$ . Let  $x \in \mathbb{R}^n$ , and define

$$A_{x,r} = \frac{1}{\operatorname{vol} B_r(x)} \int_{B_r(x)} u.$$
(7)

Campanato's Lemma is then

Lemma 2.1 (Campanato's Lemma) If there is a constant c such that

$$\int_{B_r(x)} (f - A_{x,r})^2 \le cr^{2\alpha + n} \tag{8}$$

for all x and for small r then  $f \in C^{\alpha}$ .

**Proof** First we will bound  $|A_{x,r} - A_{x,2r}|$  Calculate

$$|A_{x,r} - A_{x,2r}|^2 = \left(\frac{1}{\operatorname{vol} B_r(x)} \int_{B_r(x)} f - \frac{1}{\operatorname{vol} B_{2r}(x)} \int_{B_{2r}(x)} f\right)^2 \tag{9}$$

$$= \left(\frac{1}{\operatorname{vol} B_r(x)} \int_{B_r(x)} \left(f - \frac{1}{\operatorname{vol} B_{2r}(x)} \int_{B_{2r}(x)} f\right)\right)^2 \tag{10}$$

$$\leq \left(\frac{1}{\operatorname{vol} B_r(x)} \int_{B_r(x)} |f - A_{x,2r}|\right)^2 \tag{11}$$

$$\leq \left(\frac{2^{n}}{\text{vol } B_{2r}(x)} \int_{B_{2r}(x)} |f - A_{x,2r}|\right)^{2}.$$
 (12)

By Cauchy-Schwarz  $(\int gh)^2 \leq (\int g^2) (\int h^2)$ . Apply this with g = 1 and  $h = f - A_{x,2r}$  to get

$$|A_{x,r} - A_{x,2r}|^2 \le \frac{2^{2n}}{\text{vol } B_{2r}(x)} \int_{B_{2r}(x)} (f - A_{x,2r})^2.$$
(13)

Now use the condition, 8, so

$$|A_{x,r} - A_{x,2r}|^2 \le \frac{2^{2n} c(2r)^{2\alpha + n}}{\text{vol } B_{2r}(x)},\tag{14}$$

and we can pick a new dimensional constant  $C = \sqrt{\frac{c2^{2n+2\alpha}}{\text{vol }B_1(x)}}$  so that

$$|A_{x,r} - A_{x,2r}| \le Cr^{\alpha}.$$
(15)

We apply this to our problem. Note that  $A_{x,2^{-k}r} - A_{x,r} = \sum_{i=0}^{i=k-1} A_{x,2^{-i-1}r} - A_{x,2^{-i}r}$ , so

$$|A_{x,2^{-k_r}} - A_{x,r}| \leq \sum_{i=0}^{i=k-1} |A_{x,2^{-i-1}r} - A_{x,2^{-i}r}|$$
(16)

$$\leq \sum_{i=0}^{i=k-1} C(2^{-i-1}r)^{\alpha}$$
(17)

$$\leq Cr^{\alpha}2^{-\alpha}\sum_{i=0}^{i=k-1}2^{-i\alpha}.$$
 (18)

This is simply the sum of a geometric series, so we can use the usual formula to get

$$|A_{x,2^{-k_r}} - A_{x,r}| \le \frac{1 - 2^{-k\alpha}}{1 - 2^{-\alpha}} C r^{\alpha} 2^{-\alpha}.$$
(19)

Now let  $k \to \infty$  to get

$$|f(x) - A_{x,r}| \le \frac{1}{1 - 2^{-\alpha}} C r^{\alpha} 2^{-\alpha}.$$
(20)

Now pick another point  $y \in \mathbb{R}^n$ . Clearly we also have  $|f(y) - A_{y,s}| \leq \frac{1}{1-2^{-\alpha}} C s^{\alpha} 2^{-\alpha}$ . Now we'll estimate  $|A_{x,r} - A_{y,s}|$ . Calculate

$$|A_{x,r} - A_{y,s}|^2 = \left| \frac{1}{\text{vol } B_r(x)} \int_{B_r(x)} f - A_{y,s} \right|^2$$
(21)

$$= \left| \frac{1}{\text{vol } B_r(x)} \int_{B_r(x)} (f - A_{y,s}) \right|^2$$
(22)

$$\leq \left(\frac{1}{\operatorname{vol} B_r(x)} \int_{B_r(x)} |f - A_{y,s}|\right)^2, \tag{23}$$

and apply Cauchy-Schwaz as before to get

$$|A_{x,r} - A_{y,s}|^2 \le \frac{1}{\text{vol } B_r(x)} \int_{B_r(x)} (f - A_{y,s})^2$$
(24)

Pick r = |x - y| and s = 2|x - y| so that  $B_r(x) \subset B_s(y)$ . Therefore

$$|A_{x,r} - A_{y,s}|^2 \leq \frac{1}{\text{vol } B_r(x)} \int_{B_s(y)} (f - A_{y,s})^2$$
(25)

$$\leq \frac{\operatorname{vol} B_s(y)}{\operatorname{vol} B_r(x)} \frac{1}{\operatorname{vol} B_s(y)} \int_{B_s(y)} (f - A_{y,s})^2$$
(26)

and, by our hypothesis,

$$|A_{x,r} - A_{y,s}|^2 \le \frac{\text{vol } B_s(y)}{\text{vol } B_r(x)} cs^{2\alpha + n}.$$
(27)

Plugging in our values for r and s we have

$$|A_{x,r} - A_{y,s}|^2 \le 2^n c(2|x-y|)^{2\alpha+n},$$
(28)

and we're only interested in small |x - y|, so we can take  $2|x - y| \le 1$  to get

$$|A_{x,r} - A_{y,s}|^2 \le 2^n c(2|x-y|)^{2\alpha}.$$
(29)

At last we can calculate

$$|f(x) - f(y)| \leq |f(x) - A_{x,r}| + |A_{x,r} - A_{y,s}| + |A_{y,s} - f(y)|$$
(30)

$$\leq \frac{1}{1-2^{-\alpha}}Cr^{\alpha}2^{-\alpha} + 2^{n/2}\sqrt{c}(2|x-y|)^{\alpha} + \frac{1}{1-2^{-\alpha}}Cs^{\alpha}2^{-\alpha}$$
(31)

$$\leq \left(\frac{C2^{-\alpha}}{1-2^{-\alpha}} + 2^{n/2+\alpha}\sqrt{c} + \frac{1}{1-2^{-\alpha}}C\right)|x-y|^{\alpha}.$$
 (32)

This completes the proof.

## 3 Morrey's Lemma

Morrey's lemma is very similar to Campanato's lemma, but uses a condition on Dirichlet energy rather than  $L_2$  norm.

**Lemma 3.1** (Morrey's lemma) Let  $f : \mathbb{R}^n \to \mathbb{R}$  be continuous. If there is a constant  $c_1$  with

$$\int_{B_r(x)} |\nabla f|^2 \le c_1 r^{n-2-2\alpha} \tag{33}$$

for all x and small r then  $f \in C^{\alpha}$ .

**Proof** This is a straightforward consequence of Campanato and Poincare. By Poincare there is a constant  $c_2$  such that

$$\int_{B_r(x)} (f - A_{x,r})^2 \le c_2 r^2 \int_{B_r(x)} |\nabla f|^2 \le c_1 c_2 r^{n+2\alpha}$$
(34)

by hypothesis. The result then follows by Campanato.