

# Convection problem set

## 1. Hele-Shaw Cell Convection

The Hele-Shaw cell has a thin layer of fluid confined between two glass or plexiglass plates at  $y = 0$  and  $y = \delta$ . It allows us to study real analogues to 2-D flows.

1) Assume that  $u$  and  $w$  have the characteristic parabolic profile in  $y$

$$u(x, y, z, t) = u(x, z, t)f(y) \quad \text{and} \quad w(x, y, z, t) = w(x, z, t)f(y) \quad , \quad f(y) = 6\frac{y}{\delta} \left(1 - \frac{y}{\delta}\right)$$

and that pressure and buoyancy are independent of  $y$  and  $v = 0$ . Take the  $y$  average of the scaled equations

$$\begin{aligned} \frac{1}{Pr} \frac{D}{Dt} \mathbf{u} &= -\nabla p + Ra b \hat{\mathbf{z}} + \nabla^2 \mathbf{u} \\ \nabla \cdot \mathbf{u} &= 0 \\ \frac{D}{Dt} b &= \mathbf{e} + \nabla^2 b \end{aligned}$$

( $b$  being the deviation from  $\bar{b} = -z$ ). Show that the averaged momentum equation reduces to

$$\frac{12}{\delta^2} \mathbf{u} = -\nabla p + Ra b \hat{\mathbf{z}}$$

when  $\delta$  is very small (thin fluid layer) and the Rayleigh number is big.

2) Define the streamfunction and the  $y$ -component of the vorticity. Eliminate the pressure and write coupled equations for  $\psi$  and  $b$ .

3) Solve the linear stability problem and find the critical Rayleigh number assuming free-slip conditions on top and bottom.

4) Consider now the nonlinear problem. We could solve for weakly supercritical conditions by assuming  $\epsilon = Ra - Ra_c$  is small and expanding in suitable powers. However, it is easier to follow Lorenz and examine a truncated systems. Derive the equivalents to the Lorenz equations for this system. You can use the same expansions for  $b$  and the vorticity as in the notes. Show that the resulting equations have a simple bifurcation from the motionless state to a stable steady state

## 2. Line Plume

A line plume is driven by a steady buoyancy flux per unit length  $G_0$  (units of  $m^3 s^{-3}$ ) on a line which extends infinitely far in one horizontal direction  $y$ . The resultant plume is therefore 2-dimensional, with no variation in the  $y$ -direction. Entrainment into the plume occurs only in the  $x$ -direction. In this problem, we'll examine the derivation of the plume equations and the solutions.

1) Let the turbulent plume extend from  $-r(z) < x < r(z)$ . If the boundary were an impermeable surface, then we would have

$$u = w \frac{\partial r}{\partial z}$$

but entrainment adds an extra flow which we'll call  $-u_e$  (negative since the entrainment is going inward)

$$u = w \frac{\partial r}{\partial z} - u_e$$

Integrate the two-D continuity equation from 0 to  $r$  and show that the averaged vertical velocity

$$W = \frac{1}{r} \int_0^r dx w$$

satisfies

$$\frac{\partial}{\partial z} W r = u_e$$

2) Now consider the flux of buoyancy, integrating from 0 to  $r$  where the outside buoyancy value is  $b_e$ . Assume you can replace the average of  $wb$  with  $W$  times  $B$ , the average of  $b$ .

3) Finally treat the vertical momentum. The additional assumption is that the pressure is the same as the environment.

4) Now look for similarity solutions: assume  $W \sim z^n$ ,  $r \sim z^m$  and find the solutions given a uniform background buoyancy  $b_e \neq 0$ . Make the entrainment approximation  $u_e = \alpha W$ , where  $\alpha$  is a constant.

5) Suppose you can assess the effects of rotation by finding the height  $z_f$  beyond which the plume Rossby number  $Ro = W/(fr)$  is less than 1 using the similarity solutions above. What is  $z_f$  and what do you think might happen at greater heights?